

Interpretation of Special Relativity in the Language of classical Kinematics

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A *translation* is a function between formulas of languages preserving the logical connectives, i.e. $Tr(\phi \wedge \psi) = Tr(\phi) \wedge Tr(\psi)$, etc.

An *interpretation* of theory \mathbf{Th}_1 in theory \mathbf{Th}_2 is a translation Tr which translates all axioms of \mathbf{Th}_1 into theorems of \mathbf{Th}_2 .

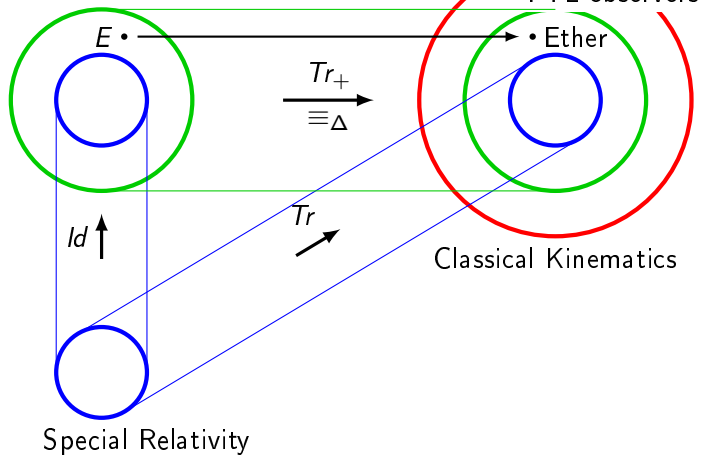
An interpretation is a *definitional equivalence* if a translated formula can be translated back such that it becomes a formula which is equivalent to the original formula.

We use these concepts to show the *differences* between theories which are *not* equivalent.

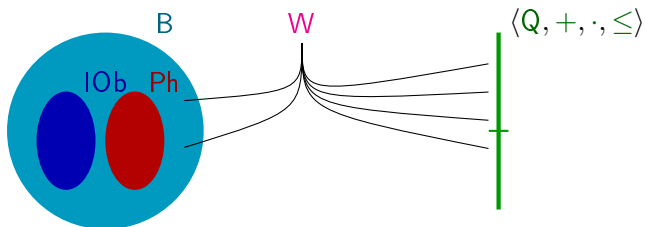
There are translations Tr , Tr_+ , and Tr_+^{-1} between the languages of NK and SR such that:

- $NK \vdash Tr(SR)$
- $NKnoFTL \vdash Tr_+(SRwithEther)$
- $SRwithEther \vdash Tr_+^{-1}(NKnoFTL)$
- Definitional equivalence: $SRwithEther \equiv_{\Delta} NKnoFTL$, i.e., Tr_+ and Tr_+^{-1} are inverses of each other up to logical equivalence in $NKnoFTL$ and $SRwithEther$.

Special Relativity with Ether



Language: $\{ B, IOb, Ph, Q, +, \cdot, \leq, W \}$



$B \iff$ Bodies (things that move)

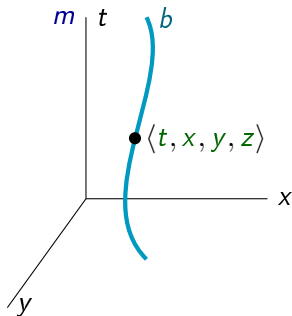
$IOb \iff$ Inertial Observers $Ph \iff$ Photons (light signals)

$Q \iff$ Quantities

$+, \cdot$ and $\leq \iff$ field operations and ordering

$W \iff$ Worldview (a 6-ary relation of type $BBQQQQ$)

$W(m, b, t, x, y, z) \iff$ “observer m coordinatizes body b at spacetime location $\langle t, x, y, z \rangle$.”



Worldline of body b according to observer m

$$wl_m(b) := \{ \langle t, x, y, z \rangle \in Q^4 : W(m, b, t, x, y, z) \}$$

$$\mathbf{Kin} := \{ \text{AxEField}, \text{AxEv}, \text{AxSelf}, \text{AxSymD}, \text{AxLine}, \text{AxTriv} \}$$
$$\mathbf{ClassicalKin}_{Full} := \mathbf{Kin} \cup \{ \text{AxEther}, \text{AbsTime}, \text{AxThExp}_{+}^{\uparrow} \}$$
$$\mathbf{SpecRel}_{Full} := \mathbf{Kin} \cup \{ \text{AxPh}_c, \text{AxThExp}^{\uparrow} \}$$

AxEField :

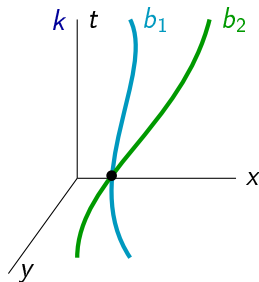
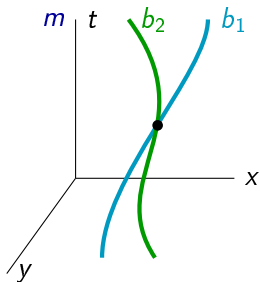
The **structure of quantities** $\langle Q, +, \cdot, \leq \rangle$ is an *Euclidean field*,

- Real numbers: \mathbb{R} ,
- Real algebraic numbers: $\overline{\mathbb{Q}} \cap \mathbb{R}$,
- Hyperreal numbers: \mathbb{R}^* ,
- Real constructable numbers,
- Etc...

AxEv :

Inertial observers coordinatize the same events (meetings of bodies).

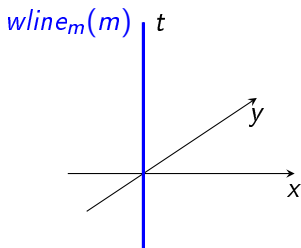
$$ev_m(\bar{x}) := \{b : W(m, b, \bar{x})\}$$



$$\forall m k \bar{x} [IOb(m) \wedge IOb(k) \rightarrow \exists \bar{y} ev_m(\bar{x}) = ev_k(\bar{y})].$$

AxSelf :

Every *Inertial observer* is stationary according to *himself*.

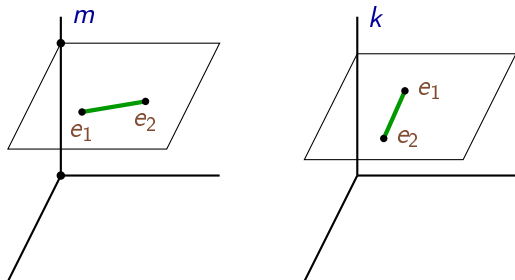


$$\forall mtxyz \left(\text{IOb}(m) \rightarrow [\text{W}(m, m, t, x, y, z) \leftrightarrow x = y = z = 0] \right).$$

AxSymD :

Inertial observers agree as to the spatial distance between two events if these two events are simultaneous for both of them.

$$\text{space}(\bar{x}, \bar{y}) := \sqrt{(x_2 - y_2)^2 + \dots + (x_d - y_d)^2}$$



$$\forall mk \bar{x} \bar{y} \bar{x}' \bar{y}' [IOb(m) \wedge IOb(k) \wedge x_1 = y_1 \wedge x'_1 = y'_1 \wedge ev_m(\bar{x}) = ev_k(\bar{x}') \wedge ev_m(\bar{y}) = ev_k(\bar{y}') \rightarrow \text{space}(\bar{x}, \bar{y}) = \text{space}(\bar{x}', \bar{y}')]]$$

AxLine :

The worldlines of inertial observers are straight lines according to inertial observers.

$$\forall mk\bar{x}\bar{y}\bar{z} \left(IOb(m) \wedge IOb(k) \wedge W(m, k, \bar{x}) \wedge W(m, k, \bar{y}) \wedge W(m, k, \bar{z}) \right. \\ \left. \rightarrow \exists a [\bar{z} - \bar{x} = a(\bar{y} - \bar{x}) \vee \bar{y} - \bar{z} = a(\bar{z} - \bar{x})] \right).$$

(In *SpecRel*, AxLine is a theorem, so including it as an axiom is redundant.)

AxTriv :

Any trivial transformation of an inertial frame is also an inertial frame.

$\forall T \in Triv[\forall m \exists k(w_{mk} = T)]$, where Triv is the set of Trivial transformations, i.e. transformations that are isometries on space and translations on time.

AxAbsTime :

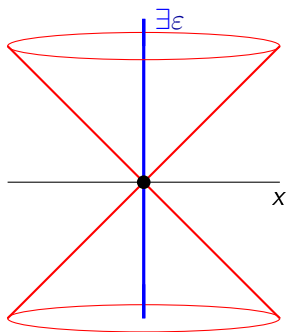
The time difference between two *events* is the same for all *inertial observers*.

$$\text{time}(\bar{x}, \bar{y}) := |x_1 - y_1|$$

$$\forall mk\bar{x}\bar{y}\bar{x}'\bar{y}' [IOb(m) \wedge IOb(k) \wedge ev_m(\bar{x}) = ev_k(\bar{x}') \wedge ev_m(\bar{y}) = ev_k(\bar{y}') \\ \rightarrow \text{time}(\bar{x}, \bar{y}) = \text{time}(\bar{x}', \bar{y}')].$$

AxEther(Einstein's AxLight) :

There exists an *inertial observer* in which the *light cones* are *right*.

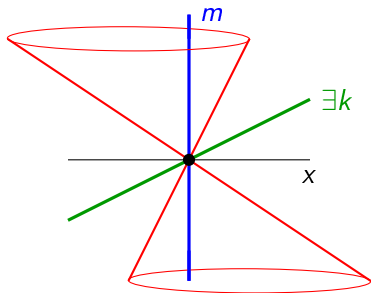


$$\exists \epsilon c \left[\text{IOb}(\epsilon) \wedge c > 0 \wedge \forall \bar{x} \bar{y} \left(\exists p \left[\text{Ph}(p) \wedge \text{W}(\epsilon, p, \bar{x}) \right. \right. \right. \\ \left. \left. \left. \wedge \text{W}(\epsilon, p, \bar{y}) \right] \leftrightarrow \text{space}(\bar{x}, \bar{y}) = c \cdot \text{time}(\bar{x}, \bar{y}) \right) \right]$$

$\text{AxThExp}_+^\uparrow$:

Inertial observers can move along any non-horizontal straight line.

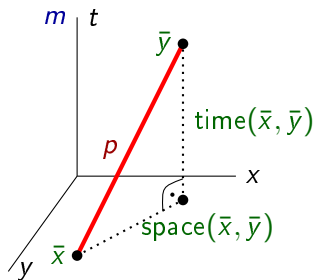
$$k \uparrow m \stackrel{\text{def}}{\iff} \text{ev}_k(\bar{x}) = \text{ev}_m(1, 0, 0, 0) \wedge \text{ev}_k(\bar{y}) = \text{ev}_m(0, 0, 0, 0) \rightarrow x_1 > y_1$$



$$\begin{aligned} \exists h (IOb(h)) \wedge \forall m \bar{x} \bar{y} (IOb(m) \wedge x_1 \neq y_1 \\ \rightarrow \exists k IOb(k) \wedge W(m, k, \bar{x}) \wedge W(m, k, \bar{y}) \wedge m \uparrow k). \end{aligned}$$

AxPh_c :

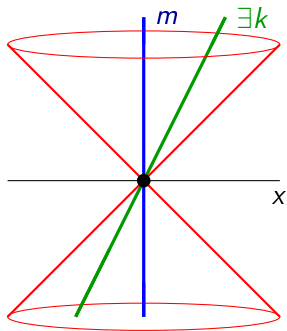
For any *inertial observer*, the *speed of light* is the same in every *direction everywhere*, and it is finite. Furthermore, it is possible to send out a *light signal* in any *direction*.



$$\exists c \left[c > 0 \wedge \forall m \bar{x} \bar{y} \left(\text{IOb}(m) \rightarrow \exists p \left[\text{Ph}(p) \wedge \text{W}(m, p, \bar{x}) \right. \right. \right. \\ \left. \left. \left. \wedge \text{W}(m, p, \bar{y}) \right] \leftrightarrow \text{space}(\bar{x}, \bar{y}) = c \cdot \text{time}(\bar{x}, \bar{y}) \right) \right]$$

AxThExp^\uparrow :

Inertial observers can move with any speed slower than that of light.



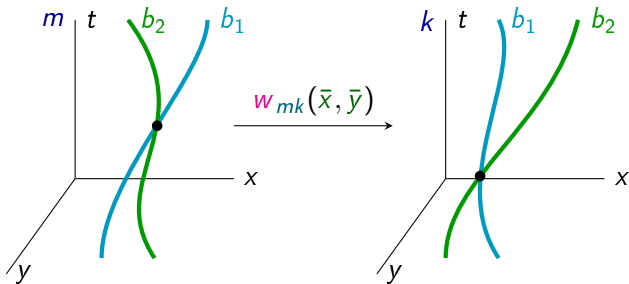
$$\begin{aligned} \exists h [IOb(h)] \wedge \forall m \bar{x} \bar{y} (IOb(m) \wedge \text{space}(\bar{x}, \bar{y}) < c \cdot \text{time}(\bar{x}, \bar{y})) \\ \rightarrow \exists k [IOb(k) \wedge W(m, k, \bar{x}) \wedge W(m, k, \bar{y}) \wedge m \uparrow k]. \end{aligned}$$

$$\mathbf{Kin} := \{ \text{AxEField}, \text{AxEv}, \text{AxSelf}, \text{AxSymD}, \text{AxLine}, \text{AxTriv} \}$$

$$\mathbf{ClassicalKin}_{Full} := \mathbf{Kin} \cup \{ \text{AxEther}, \text{AbsTime}, \text{AxThExp}_{+}^{\uparrow} \}$$

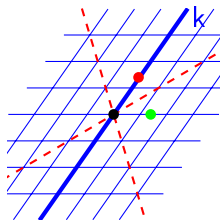
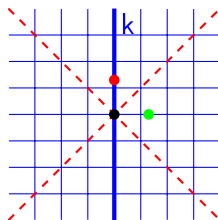
$$\mathbf{SpecRel}_{Full} := \mathbf{Kin} \cup \{ \text{AxPh}_c, \text{AxThExp}^{\uparrow} \}$$

Worldview transformation :

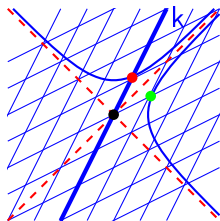
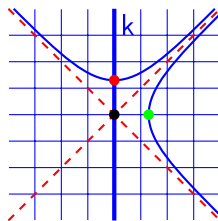


$$w_{mk}(\bar{x}, \bar{y}) \stackrel{\text{def}}{\iff} ev_m(\bar{x}) = ev_k(\bar{y})$$

Galilean transformations:

worldview of m worldview of k

Poincaré transformations:

worldview of m worldview of k

Representation Theorems:

Theorem:

$\text{SpecRel}_{\text{Full}} \vdash \forall mk [IOb(m) \wedge IOb(k) \rightarrow \text{"}w_{mk} \text{ is a Poincaré Transformation"}]$.

- ca. 1998 proven for a version of BasAx strongly related to SpecRel in the "Big Book" by H. Andréka, J. X. Madarász & I. Németi
- Synthese 2012 "A logic road from special relativity to general relativity" by H. Andréka, J. X. Madarász, I. Németi & G. Székely (Theorem 2.2)

Theorem:

$\text{ClassicalKin}_{\text{Full}} \vdash \forall mk [IOb(m) \wedge IOb(k) \rightarrow \text{"}w_{mk} \text{ is a Galilean Transformation"}]$.

- proof for a different version of classical Kinematics in the "Big Book" by H. Andréka, J. X. Madarász & I. Németi

Theorem:

$\text{SpecRel}_{\text{Full}} \vdash \forall k P[\text{IOb}(k) \wedge \text{"}P \text{ is a (orthochronous) Poincaré Transformation"} \rightarrow \exists k' w_{kk'} = P].$

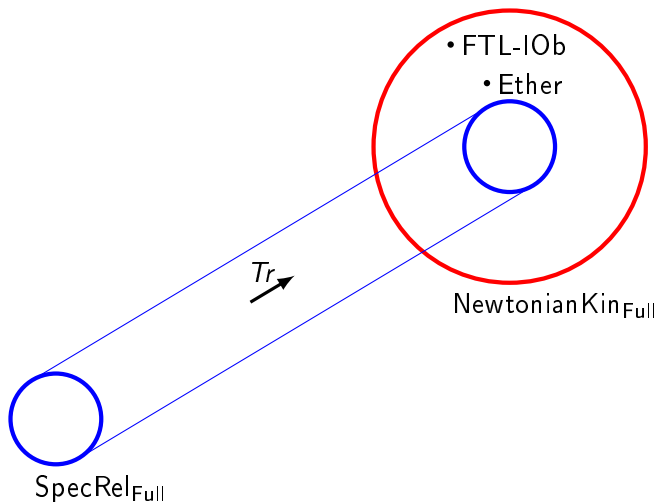
Theorem:

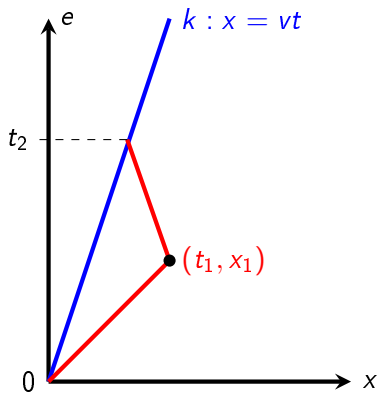
$\text{ClassicalKin}_{\text{Full}} \vdash \forall k G[\text{IOb}(k) \wedge \text{"}G \text{ is a (orthochronous) Galilean Transformation"} \rightarrow \exists k' w_{kk'} = G].$

- can be proven based on the ideas in the "Big Book" by H. Andréka, J. X. Madarász & I. Németi

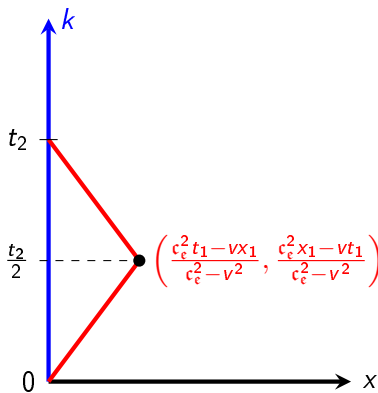
Theorem:

There is an interpretation Tr of SpecRel_{Full} in $\text{ClassicalKin}_{Full}$.

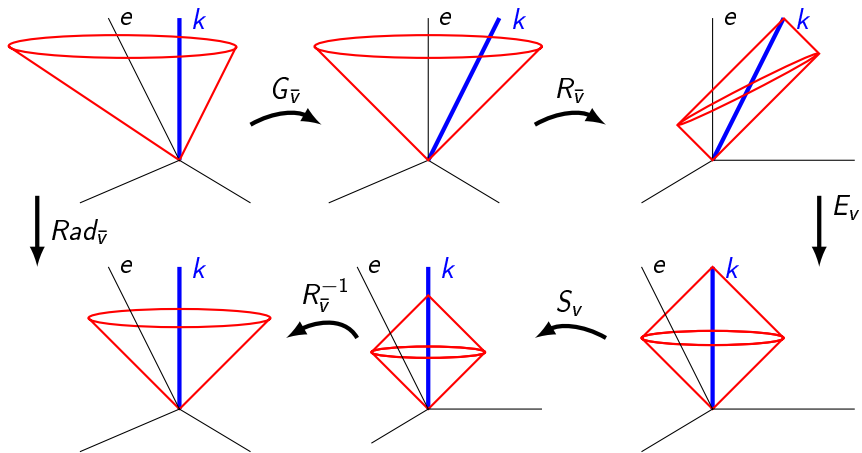




$$A = \begin{bmatrix} \frac{1}{1-v^2} & \frac{-v}{1-v^2} & 0 & 0 \\ \frac{-v}{1-v^2} & \frac{1}{1-v^2} & 0 & 0 \\ 0 & 0 & \frac{1}{\sqrt{1-v^2}} & 0 \\ 0 & 0 & 0 & \frac{1}{\sqrt{1-v^2}} \end{bmatrix}$$



$$S = \sqrt{1 - v^2} \quad (\text{if } c = 1)$$



$$Rad_k = R^{-1} \circ S \circ A \circ R \circ G$$

$$\text{Tr}(a + b = c) := a + b = c$$

$$\text{Tr}(a \cdot b = c) := a \cdot b = c$$

$$\text{Tr}(a < b) := a < b$$

$$\text{Tr}(\mathbf{W}_{SR}(k, b, t, x, y, z)) :=$$

$$\exists t'x'y'z'[\mathbf{W}_{NK}(k, b, t', x', y', z') \wedge \text{Rad}_k(t', x', y', z') = (t, x, y, z)]$$

$$\text{Tr}(\mathbf{IOb}_{SR}(k)) := \mathbf{IOb}_{NK}(k) \wedge \forall \varepsilon [\text{Ether}(\varepsilon) \rightarrow \text{speed}_\varepsilon^{NK}(k) < c]$$

where

$$\text{Ether}(\varepsilon) \stackrel{\text{def}}{\iff} \mathbf{IOb}_{NK}(\varepsilon) \wedge \forall p [\text{Ph}(p) \rightarrow \text{speed}_\varepsilon^{NK}(p) = c]$$

and

$$\text{speed}_m(k) = v \stackrel{\text{def}}{\iff} \forall \bar{x}, \bar{y} \in \mathbb{Q}^d [\mathbf{W}(m, k, \bar{x}) \wedge \mathbf{W}(m, k, \bar{y}) \wedge x_1 \neq y_1 \rightarrow \frac{\text{space}(\bar{x}, \bar{y})}{\text{time}(\bar{x}, \bar{y})} = v]$$

Example: Translation of AxSelf_{SR}

$$\begin{aligned} & \text{Tr}[\forall m \text{IOb}_{SR}(m) \rightarrow \forall \bar{x} (\text{W}_{SR}(m, m, \bar{x}) \leftrightarrow x_2 = \dots = x_d = 0)] \\ & \equiv \forall m [\text{IOb}_{NK}(m) \wedge \forall \varepsilon [\text{Ether}(\varepsilon) \rightarrow \text{speed}_{\varepsilon}^{NK}(m) < c] \\ & \rightarrow \forall \bar{x} (\exists \bar{y} [\text{W}_{NK}(m, m, \bar{y}) \wedge \text{Rad}_k(\bar{y}) = \bar{x}] \leftrightarrow x_2 = \dots = x_d = 0)] \end{aligned}$$

Theorem:

There is an interpretation Tr of $\mathbf{SpecRel}_{Full}$ in $\mathbf{ClassicalKin}_{Full}$.
 $\mathbf{ClassicalKin}_{Full} \vdash Tr(\mathbf{SpecRel}_{Full})$

- $\mathbf{ClassicalKin}_{Full} \vdash Tr(\mathbf{AxEField}_{SR})$
- $\mathbf{ClassicalKin}_{Full} \vdash Tr(\mathbf{AxEv}_{SR})$
- $\mathbf{ClassicalKin}_{Full} \vdash Tr(\mathbf{AxSelf}_{SR})$
- $\mathbf{ClassicalKin}_{Full} \vdash Tr(\mathbf{AxSymD}_{SR})$
- $\mathbf{ClassicalKin}_{Full} \vdash Tr(\mathbf{AxLine}_{SR})$
- $\mathbf{ClassicalKin}_{Full} \vdash Tr(\mathbf{AxTriv}_{SR})$
- $\mathbf{ClassicalKin}_{Full} \vdash Tr(\mathbf{AxThExp}_{SR}^{\uparrow})$

ClassicalKin^{NoFTL}_{Full} := Kin \cup {AxEther, AbsTime, AxThExp[↑]_{NoFTL}, AxNoFTL}

AxThExp[↑]_{NoFTL} :

Inertial observers can move with any speed which is in the ether frame slower than that of light.

$$\begin{aligned} \exists h (IOb(h)) \wedge \forall \varepsilon \bar{x} \bar{y} (Ether(\varepsilon) \wedge \text{space}(\bar{x}, \bar{y}) < c \cdot \text{time}(\bar{x}, \bar{y})) \\ \rightarrow \exists k IOb(k) \wedge W(\varepsilon, k, \bar{x}) \wedge W(\varepsilon, k, \bar{y}) \wedge m \uparrow k). \end{aligned}$$

AxNoFTL :

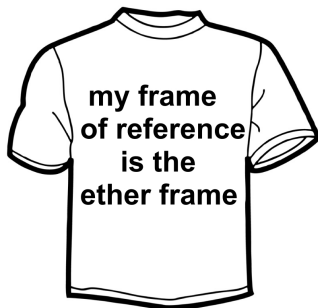
All inertial observers move slower than light with respect to the ether frames.

$$\forall m \varepsilon [IOb(m) \wedge Ether(\varepsilon) \rightarrow \text{Speed}_{\varepsilon}^{NK}(m) < c].$$

$$\text{SpecRel}_{Full}^{\varepsilon} := \text{SpecRel}_{Full} \cup \{\text{AxPrimitiveEther}\}$$

AxPrimitiveEther :

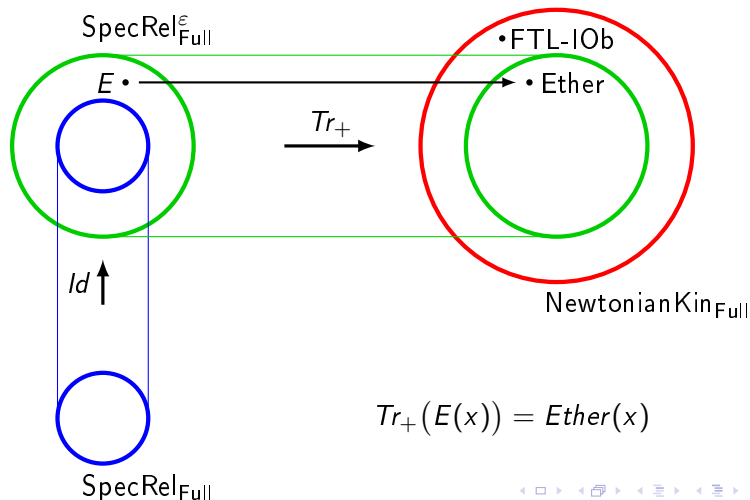
There is a non-empty class of ether observers, stationary with respect to each other, which is closed under trivial transformations.



$$\exists \varepsilon (E(\varepsilon) \wedge \forall k [[IOb(k) \wedge (\exists T \in Triv) w_{\varepsilon k}^{SR} = T] \leftrightarrow E(k)])$$

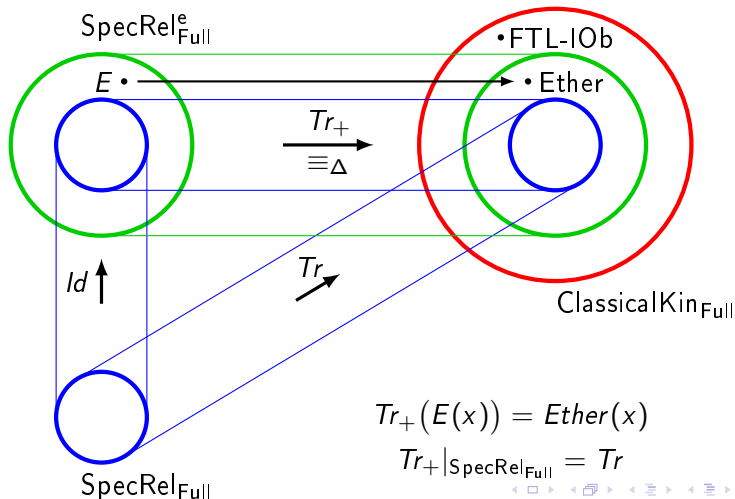
Theorem:

Tr_+ is a definitional equivalence between SpecRel_{Full}^E and $\text{ClassicalKin}_{Full}^{NoFTL}$.

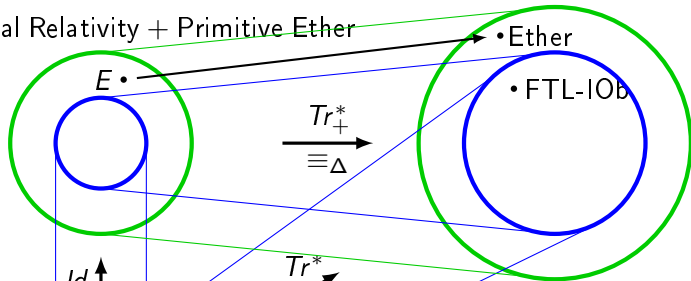


Theorem:

Tr_+ is a definitional equivalence between $SpecRel_{Full}^e$ and $ClassicalKin_{Full}^{NoFTL}$.



Special Relativity + Primitive Ether



Classical Kinematics

Special Relativity

$$v \rightarrow \frac{v}{c-v}$$